[CONTRIBUTION FROM THE CHEMICAL LABORATORY OF THE UNIVERSITY OF CALIFORNIA]

## The Heat Capacity of Nickel from 15 to 300°K. Entropy and Free Energy Functions<sup>1</sup>

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The heat capacity of nickel has been measured from 15 to 300°K. The entropy, heat content and free energy functions have been calculated.

The work reported here was done because the free energy function for nickel was needed in connection with an equilibrium study of the reduction of nickel chloride by hydrogen at temperatures up to some 700°K.

A sample of very pure nickel was prepared through the courtesy of the research laboratory of the International Nickel Company, which supplied the following analysis:

| С   | Al  | в   | Ca  | Co      | Cu      | Fe  | Рb  |
|---|-----|-----|-----|---------|---------|-----|-----|
| 0.014%                                    | VST | VST | VST | 0.0018% | 0.0009% | VST | NF  |
| Mg  | Sn  | As  | Ag  | Na      | v       | Zn  | Si  |
| VST                                       | NF  | NF  | NF  | NF      | NF      | NF  | VST |
| NF = none found. VST = very slight trace, |     |     |     |         |         |     |     |

The carbon was determined by combustion and the cobalt and copper by wet analysis. The other elements were evaluated by spectrographic comparison. By this method copper was given as VST which apparently can be taken as of the order of 0.001%. The sample had been prepared in the form of a cylinder 4.2 cm. diam.  $\times$  10.8 cm. long. The cast cylinder had been annealed at 900° for 2 hours in an atmosphere of hydrogen. Before the nickel was placed in the calorimeter the sample was heated inductively in vacuo to remove any hydrogen which may have dissolved in the nickel. It was held between 1000 and 1100° with a vacuum down to 2  $\times$  $10^{-5}$  mm. for 5 hr. and between 1100 and 1200° and a vacuum down to  $8 \times 10^{-6}$  mm. for an additional 5 hr. It was cooled in vacuo to 800° in one hour and then slowly to  $100^{\circ}$  over a period of 17 hr. before removal. The sample weighed 1309.79 g. in vacuo. The atomic weight was taken as 58.69.

The measurements were made in the calorimeter described by Giauque and Meads<sup>2</sup> except that a standard copper-constantan thermocouple with the laboratory designation W-26 was used. The thermocouple had been compared with a helium gas thermometer and was checked against the following fixed points at the conclusion of the present measurements: triple point  $(13.92^{\circ}K.)$  and boiling point  $(20.36^{\circ}K.)$  of hydrogen, and the triple point of nitrogen  $(63.15^{\circ}K.)$ . One calorie was taken as 4.1840 absolute joules.

The observed data are given in Table I.

The heat capacity of nickel has been measured by Rodebush and Michalek<sup>3</sup> from 67 to 106°K. and from 196 to 203°K.; by Simon and Ruhemann<sup>4</sup> from 71 to 83°K.; by Eucken and Werth<sup>5</sup> from 15

(1) This work was supported in part by the Office of Naval Research, United States Navy.

(2) W. F. Giauque and P. F. Meads, THIS JOURNAL, 63, 1897 (1941).

(3) W. H. Rodebush and J. C. Michalek, *ibid.*, 47, 2117 (1925).

(4) F. Simon and M. Ruhemann, Z. physik. Chem., 129, 321 (1927).
(5) A. Eucken and H. Werth, Z. anorg. allgem. Chem., 188, 152

(5) A. Eucken and H. Werth, Z. anorg. allgem. Chem., 188, 152 (1930).

| <b>TABLE I</b> |
|----------------|
|----------------|

HEAT CAPACITY OF NICKEL

| $^{\circ}C. = 27$       | 3.16°K.; | cal. deg                | (g. atom) | <sup>-1</sup> ; at. wt. | = 58.69 |
|-------------------------|----------|-------------------------|-----------|-------------------------|---------|
| <i>Τ</i> , ° <b>Κ</b> . | $C_p$    | <i>Τ</i> , ° <b>Κ</b> . | $C_p$     | <i>T</i> , °K.          | $C_p$   |
| Seri                    | es I     | 82.22                   | 2.534     | 196.03                  | 5,322   |
| 12.95                   | 0.036    | 88.52                   | 2.800     | 204.43                  | 5.423   |
| 14.68                   | .037     | 95.13                   | 3.067     | 212.96                  | 5.513   |
| 16.71                   | .052     | 102.07                  | 3.335     | 221.33                  | 5.585   |
| 19.46                   | .075     | 109.01                  | 3.580     | 229.59                  | 5.655   |
| 22.78                   | .107     | 116.23                  | 3.797     | 238.59                  | 5.729   |
| 26.37                   | .159     | 123.62                  | 3.996     | 247.96                  | 5.820   |
| 30.14                   | .236     | 131.77                  | 4.205     | Series                  | II      |
| 33.83                   | .332     | 139.64                  | 4.395     | 028 75                  | 5 729   |
| 37.86                   | .458     | 147.50                  | 4.565     | 200.10                  | 5 818   |
| 42.93                   | .647     | 155.38                  | 4.722     | 256 73                  | 5 800   |
| <b>48</b> .93           | .902     | 163.33                  | 4.857     | 200.75                  | 5.060   |
| 54.89                   | 1.199    | 171.37                  | 4.989     | 200,09<br>974 Q1        | 6 021   |
| 62.12                   | 1.552    | 179.48                  | 5.112     | 214.01                  | 6 101   |
| 69.74                   | 1.932    | 187.68                  | 5.221     | 201.04                  | 6 905   |
| 75.91                   | 2.236    |                         |           | 202.90                  | 6 108   |

to 204°K.; by Bronson and Wilson<sup>6</sup> from 203 to  $303^{\circ}$ K. and by Aoyama and Kanda<sup>7</sup> from 80 to  $273^{\circ}$ K.

The results of Rodebush and Michalek are about 1.3% high in the range 67 to  $106^{\circ}$ K. and about 2.8% high near 200°K. The few measurements of Simon and Ruhemann are about 3% above the present results. The data of Eucken and Werth average about 2% low in the range 15 to  $35^{\circ}$ K., 1% high from 35 to 80°K., 0.4% low from 80 to 130°K. and 0.4% high from 130 to 200°K. The results of Bronson and Wilson are 0.4% high from 203 to 273°K. and 0.25% low from 273 to 303°K. The data of Aoyama and Kanda average about 1%high over the range 80 to  $273^{\circ}$ K. It is difficult to say how many of the above differences are due to impure samples and unknown physical state. Rodebush and Michalek used 99+% Ni. Eucken and Wilson had 0.5% Mn in their sample. Bron-son and Wilson used nickel with 0.03% Cu, 0.13%Fe, 0.10% Si and 0.03% C. Aoyama and Kanda used very pure nickel.

Keesom and Clark<sup>8</sup> have measured the heat capacity of nickel from 1 to  $20^{\circ}$ K. They give the equation

#### $C = 6.594 \times 10^{-6} T^3 + 0.001744 T$

over the above range. Their sample contained 99.81% Ni, 0.017% Cu, 0.083% Fe, 0.037% Mg and 0.04% C.

Clusius and Goldmann<sup>9</sup> have covered the range 10 to 26°K. No analysis of their sample is given.

(6) H. L. Bronson and A. J. C. Wilson, Can. J. Research, 14A, 181 (1936).

(7) S. Aoyama and E. Kanda, J. Chem. Soc. Japan, 62, 312 (1941).

(8) W. H. Keesom and C. W. Clark, Physica, 2, 513 (1935).

(9) K. Clusius and J. Goldmann, Z. physik. Chem., 31, 256 (1936)

The results of Clusius and Goldmann agree well with those of Keesom and Clark, and join well with the results of the present research. The results of Keesom and Clark made extrapolation necessary only below 1°K. in obtaining the entropy.

The thermodynamic properties of nickel calculated from the present results are given in Table II.

Calculations were extended to 800°K. on the basis of available data; however, the increments above 298.16°K. were in substantial agreement with calculations of Kelley,<sup>10</sup> who has summarized all available results as increments above 298.16°. To combine the results

$$S = 7.14 + (S - S_{298.16})_{\text{Kelley}}$$
$$\frac{H^{\circ} - H_{0}^{\circ}}{T} = \frac{3.837 \times 298.16 + (H_{T} - H_{295.16})_{\text{Kelley}}}{T}$$

and

$$\frac{F^{\circ} = H_0^{\circ}}{T} = \frac{H^{\circ} - H_0^{\circ}}{T} - S$$

#### Table II

THERMODYNAMIC PROPERTIES OF NICKEL

|                         |       |       | $H^{\circ} - H^{\circ}_{0}$ | $(F^{\circ} - H_{\theta}^{\circ})$ |
|-------------------------|-------|-------|-----------------------------|------------------------------------|
| <i>Τ</i> , ° <b>Κ</b> . | Cp    | S     | $\overline{T}$              |                                    |
| 15                      | 0.043 | 0.033 | 0.018                       | 0.015                              |
| 20                      | .081  | .050  | .029                        | .021                               |
| 25                      | .142  | ,074  | .045                        | .029                               |
| 30                      | .234  | .107  | .068                        | .039                               |
| 35                      | .367  | .153  | .101                        | .052                               |
| 40                      | .534  | .213  | .145                        | .068                               |
| 45                      | .732  | .287  | .199                        | .088                               |
| 50                      | .956  | .375  | .263                        | .112                               |

(10) K. K. Kelley, U. S. Bureau of Mines Bulletin, No. 476, 121 (1949).

| 60     | 1.448 | .592           | .418  | .174          |
|--------|-------|----------------|-------|---------------|
| 70     | 1.953 | .853           | .601  | .252          |
| 80     | 2.433 | 1.145          | ,800  | .345          |
| 90     | 2.868 | 1.456          | 1.005 | .451          |
| 100    | 3.258 | 1.778          | 1.211 | .568          |
| 110    | 3.605 | 2.105          | 1.413 | . 69 <b>2</b> |
| 120    | 3.903 | 2,432          | 1,608 | .824          |
| 130    | 4.164 | 2,755          | 1.795 | .960          |
| 140    | 4.401 | 3.073          | 1.974 | 1.099         |
| 150    | 4.614 | 3,384          | 2.143 | 1.241         |
| 160    | 4.802 | 3.688          | 2.304 | 1.384         |
| 170    | 4.969 | 3.984          | 2.455 | 1.529         |
| 180    | 5.118 | 4.272          | 2.599 | 1.673         |
| 190    | 5.250 | 4.552          | 2.735 | 1.817         |
| 200    | 5.370 | 4.824          | 2.863 | 1.961         |
| 210    | 5.477 | 5.089          | 2,986 | 2.103         |
| 220    | 5.573 | 5.346          | 3.101 | 2.245         |
| 230    | 5.662 | 5. <b>5</b> 96 | 3.211 | 2.385         |
| 240    | 5.747 | 5.839          | 3.315 | 2.524         |
| 250    | 5.831 | 6.075          | 3.413 | 2.662         |
| 260    | 5.914 | 6.305          | 3.508 | 2.797         |
| 270    | 5.996 | <b>6.53</b> 0  | 3.599 | 2.931         |
| 280    | 6.078 | 6.750          | 3.686 | 3.064         |
| 290    | 6.160 | 6.965          | 3.770 | 3.195         |
| 298.16 | 6.227 | 7.137          | 3.837 | <b>3.30</b> 0 |
| 300    | 6.242 | 7.175          | 3.851 | 3.324         |
|        |       |                |       |               |

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# The Heat Capacities and Heat Contents of Solutions of Cerium and Neodymium Chlorides at $25^{\circ}$

By F. H. Spedding and Carl F. Miller

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The change in heat capacities of solution and dilution and the change in heat content of solution and dilution of the chlorides of cerium and neodymium have been measured for concentrations up to 0.35 and 0.40 molal, respectively. Both properties of these two electrolytes show agreement with the limiting Debye-Hückel equations within limits of experimental error.

### Introduction

Previous studies on the solutions of rare earth compounds<sup>2</sup> have illustrated the need for further data on the physical properties of the rare earth elements and their compounds in regard to the role that they can play in checking physical chemistry theories.

The partial molal heat contents and partial molal heat capacities of higher charged electrolytes in solution should be sensitive to deviations from the ideal ionic solution and to deviations from the

(1) Work was performed in the Ames Laboratory of the Atomic Energy Commission.

(2) F. H. Spedding, P. E. Porter and J. M. Wright, THIS JOURNAL, 74, 2778 (1952).

theoretical equations which describe these two properties due to such phenomena as association, changes in hydration of the ions, and compound formation. Since the charge on the ion is a constant factor for all the rare earth elements and the ionic radius of these elements is known to decrease with atomic number<sup>3</sup> the heat of solution of a given type of compound of these elements should increase with atomic number. Smaller ions of the same charge would, by presenting a relatively larger surface charge density to the surrounding water molecules, orient these water dipoles more strongly and to a greater degree than larger ions of like charge.

(3) W. H. Zacharinsen, Phys. Rev., 73, 1104 (1948).